

## Spherical Harmonics

gravity field (e.g.  $\Delta p$ ,  $\Delta g$ ,  $\Delta N$ )

magnetic field

free oscillations

variations in seismic velocities

convection, ...



Introduce using gravity eqn

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial V}{\partial \theta} \right)$$

$$+ \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} = 0$$

Solve using separation of variables

$$V = R(r) \Theta(\theta) \Phi(\phi)$$

$$\text{choose } R(r) = r^l \text{ or } r^{-(l+1)}$$

$$l(l+1) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \Theta}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \Phi(\phi) \frac{\partial^2 \Phi}{\partial \phi^2} = 0$$

Note - only 3rd term has  $\phi \Rightarrow = \text{const}$

$$\frac{d^2 \Phi}{d \phi^2} = k \Phi \Rightarrow \Phi = A \cos m\phi + B \sin m\phi$$

is a soln ( $k = -m^2$ )

then  $\Theta$  eqn becomes Legendre's associated eqn

$$\frac{1}{\sin \theta} \frac{d}{d \theta} \left( \sin \theta \frac{d \Theta}{d \theta} \right) + \left[ l(l+1) - \frac{m^2}{\sin^2 \theta} \right] \Theta = 0$$

The solutions, for  $m \leq l$  is given in terms of Associated Legendre Polynomials

$$\textcircled{1} \quad P_e^m(\cos\theta) = \frac{(1-\cos^2\theta)^{m/2}}{2^l l!} \frac{d^{lm}}{d(\cos\theta)^{lm}} (\cos^2\theta - 1)^l$$

\* Finally

$$V_l(r, \theta, \phi) = \begin{bmatrix} r^l \\ -l(l+1) \end{bmatrix} P_e^m(\cos\theta) \left[ A_e^m \cos m\phi + B_e^m \sin m\phi \right]$$

$$* V = \sum_{l=0}^{\infty} \sum_{m=0}^l V_l^m$$

$l = \text{degree}$   
 $m = \text{order}$

For geophysical purposes,  $P_e^m$ 's are best normalized to have  $\text{rms} = 1$

In this form

$$P_0^0 = 1$$

$$P_3^0 = \frac{\sqrt{7}}{2} (5\cos^3\theta - 3\cos\theta)$$

$$P_1^0 = \sqrt{3} \cos\theta$$

$$P_3^1 = \sqrt{\frac{21}{8}} (5\cos^2\theta - 1) \sin\theta$$

$$P_1^1 = \sqrt{3} \sin\theta$$

$$P_3^2 = \sqrt{\frac{105}{2}} \cos\theta \sin^2\theta$$

$$P_2^0 = \frac{\sqrt{5}}{2} (3\cos^2\theta - 1)$$

$$P_3^3 = \sqrt{\frac{35}{8}} \sin^3\theta$$

$$P_2^1 = \sqrt{\frac{5}{3}} \cos\theta \sin\theta$$

$$P_2^2 = \sqrt{\frac{15}{2}} \sin^2\theta$$

The total function needs both  $\theta$  &  $\phi$  dependence

$$Y_{lm}^C = P_l^m(\cos\theta) \cos m\phi$$

$$Y_{lm}^S = P_l^m(\cos\theta) \sin m\phi$$

where

$l$  = degree

$m=0$  zonal

$m$  = order

$m=l$  sectoral

$0 < m < l$  tesseral

$l$  even  $\Rightarrow$  symmetric

$l$  odd  $\Rightarrow$  anti-symmetric

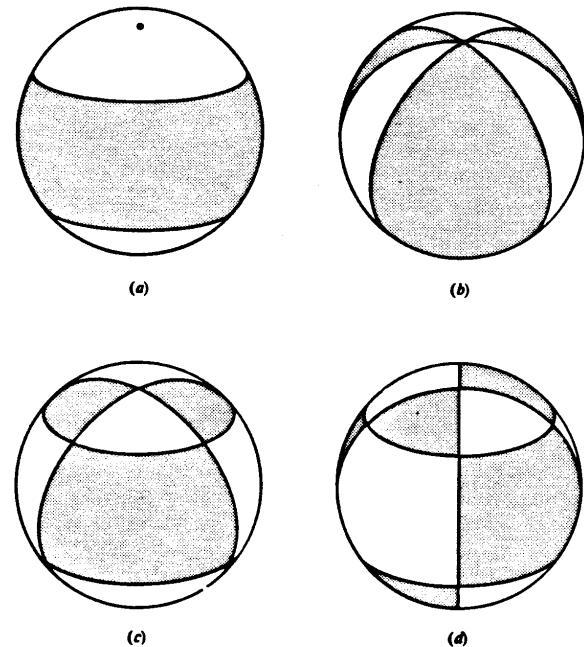
$l-m$  zero crossings in colatitude

$l-m+1$  patches in colatitude

$2m$  zero crossings in longitude

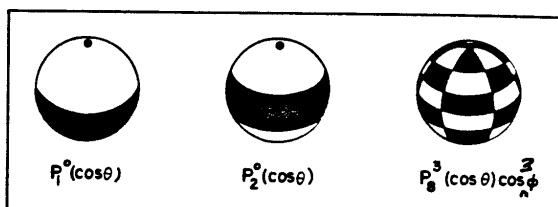
$2m$  patches in longitude

$$\text{wavelength } \lambda = \frac{2\pi a}{\sqrt{l(l+1)}} \sim \frac{2\pi a}{l + 1/2}$$

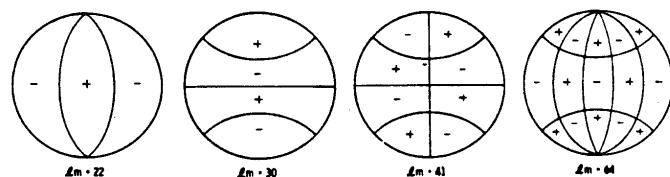


**Fig. 2.2. Spherical harmonics.** Hatched and blank regions show different signs.

- (a)  $P_0^0 = \frac{1}{4}(1 + 3 \cos 2\varphi)$ ,  
 (b)  $P_3^3 \cos 3\varphi = 15 \sin^3 \vartheta \cos 3\varphi$ ,  
 (c)  $P_4^2 \cos 2\varphi = \frac{15}{16}(3 + 4 \cos 2\vartheta - 7 \cos 4\varphi) \cos 2\varphi$ ,  
 (d)  $P_4^2 \sin 2\varphi = \frac{15}{16}(3 + 4 \cos 2\vartheta - 7 \cos 4\varphi) \sin 2\varphi$ .



**Figure A.2** Examples of surface spherical harmonics.  $P_1^0(\cos \theta)$  and  $P_2^0(\cos \theta)$  are dominant terms in the earth's magnetic and gravitational fields respectively.  $P_3^1(\cos \theta \cos \phi)$  shows a tesseral harmonic.



**Figure C.1.** Examples of spherical harmonics.  $m = 0$  gives zonal harmonics,  $m = l$  gives sectoral harmonics and the general case  $0 < m < l$  are known as tesseral harmonics. Reproduced, by permission, from Kaula (1968).

Expression of a scalar field in spherical harmonics

$$f(\theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=0}^l P_l^m(\cos \theta) [a_l^m \cos m\phi + b_l^m \sin m\phi]$$

Examples: Topography  
Gravity  
Magnetic field

Spectrum: root mean square coefficient

$$S_l^{\text{rms}} = \left( \frac{\sum_{m=0}^l (a_l^m)^2 + (b_l^m)^2}{(2l+1)} \right)^{1/2}$$

Vector spherical harmonics - e.g. displacement

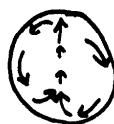
2 types

Polaroidal (spheroidal)

$$u_r = U_r(r) P_l^m(\cos \theta) \begin{bmatrix} \cos m\phi \\ \sin m\phi \end{bmatrix}$$

$$u_\theta = V_\theta(r) \frac{\partial P_l^m}{\partial \theta} \begin{bmatrix} \cos m\phi \\ \sin m\phi \end{bmatrix}$$

$$u_\phi = V_\phi(r) P_l^m \frac{\partial}{\partial \phi} \begin{bmatrix} \cos m\phi \\ \sin m\phi \end{bmatrix}$$



(section)

Toroidal



$$u_\theta = \frac{W(r)}{\sin\theta} P_e^m \frac{\partial}{\partial\theta} \begin{cases} \cos m\phi \\ \sin m\phi \end{cases}$$

$$u_\phi = -W(r) \frac{\partial P_e^m}{\partial\theta} \begin{cases} \cos m\phi \\ \sin m\phi \end{cases}$$

Spectrum

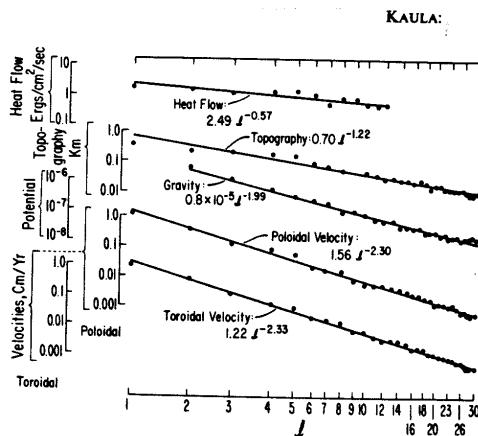


Fig. 1. Spectral magnitudes of observed surface fields

Stacy

8.1 THE MAIN FIELD 217

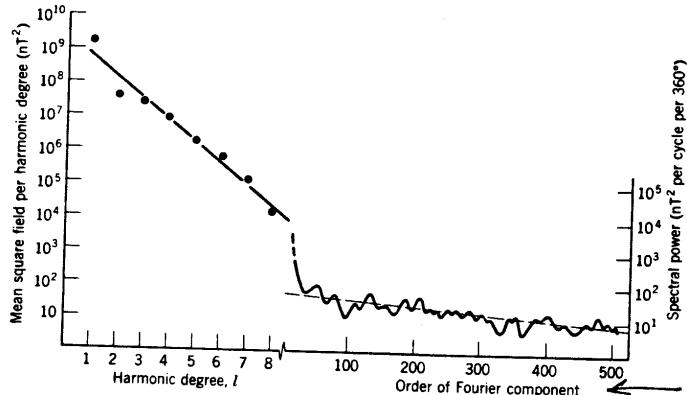
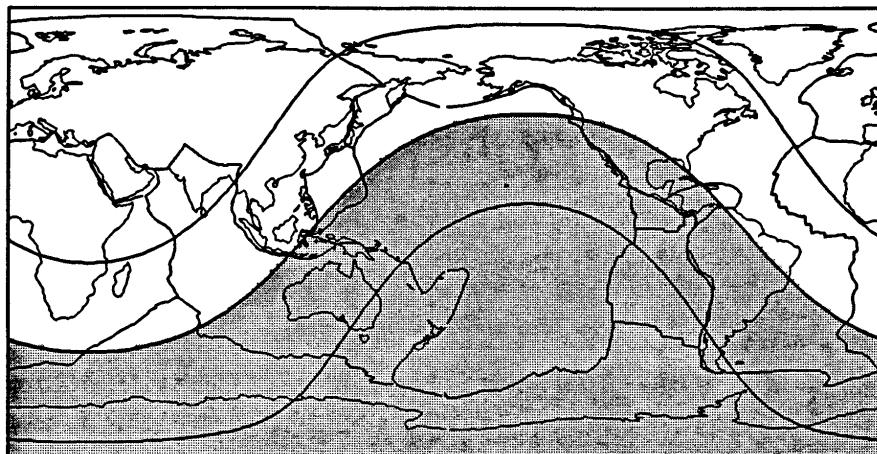


Figure 8.2. Spatial spectrum of the main geomagnetic field. The left part of the figure gives the mean square amplitudes of the field represented by each harmonic degree, for low degrees, according to Lowes (1974) and the right-hand part gives the higher-order terms of a Fourier analysis by Bullard (1967), corrected by Lowes (1974), of a world-encircling line profile by Alldredge et al. (1963).

asymptotic limit

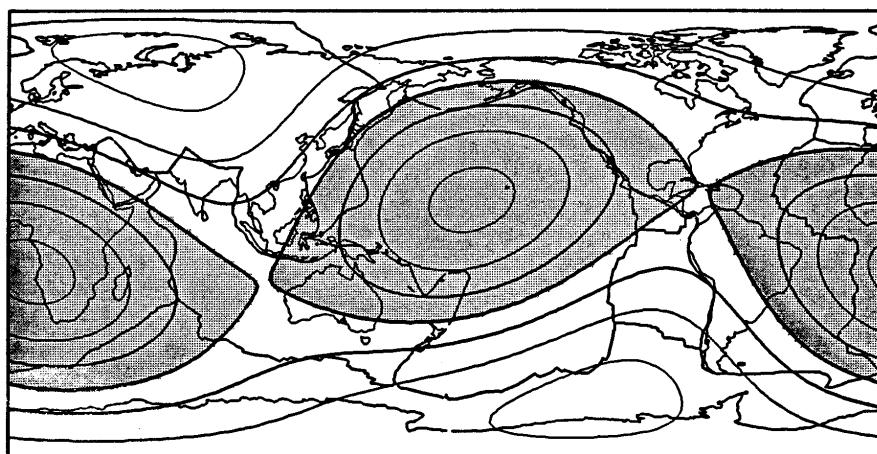
asymptotic limit

Global Topography: degree 1-1



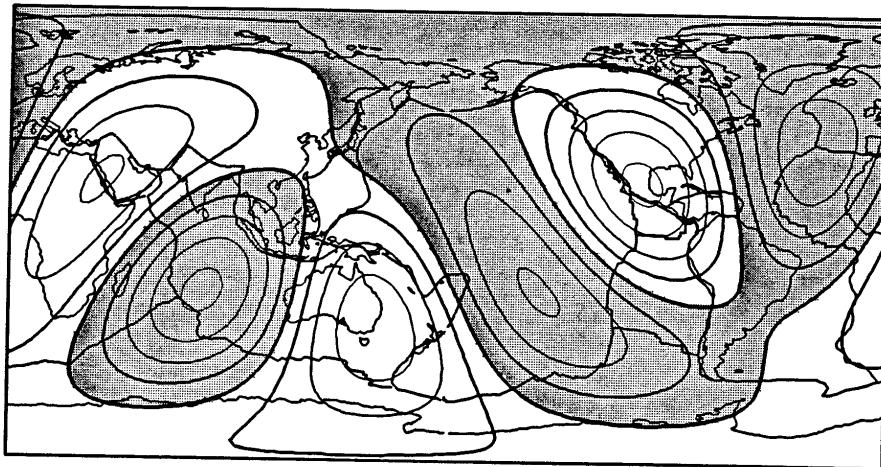
contour interval: 1000 m ; lows shaded

Global Topography: degree 2-2



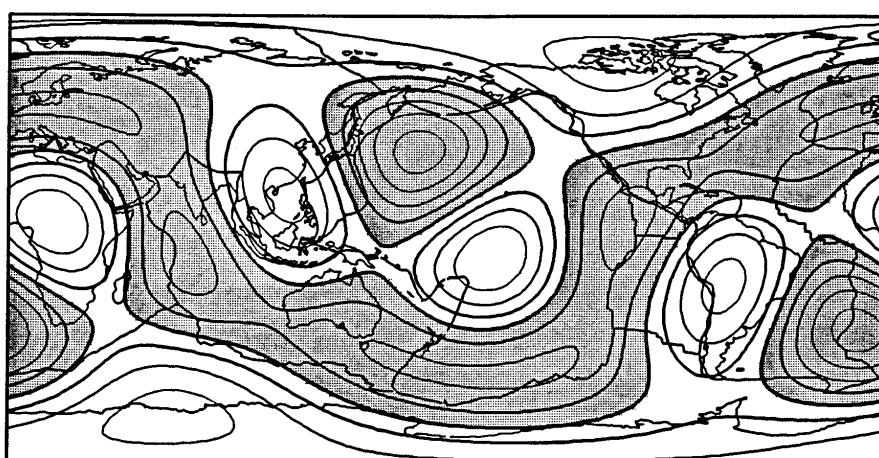
contour interval: 500 m

Global Topography: degree 3-3



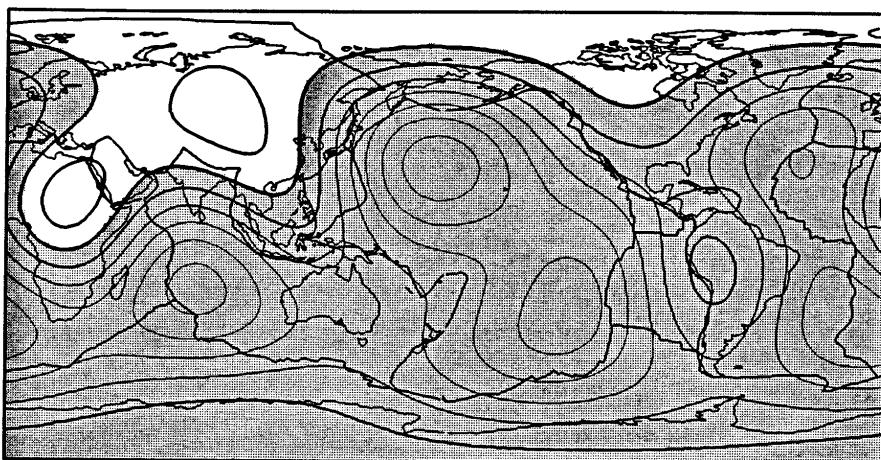
contour interval: 500 m

Global Topography: degree 4-4



contour interval: 500 m

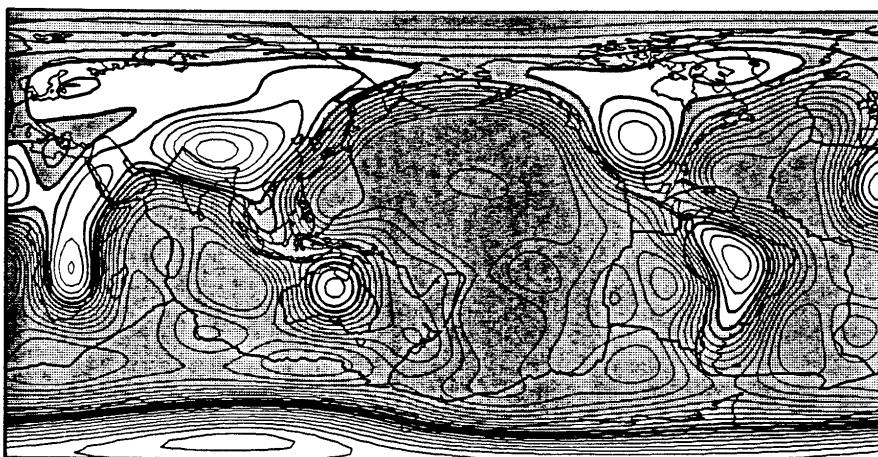
Global Topography: degree 0-4



contour interval: 1000 m

10

Global Topography: degree 0-10



contour interval: 500 m

Global Topography: degree 0-20



contour interval: 1000 m

Relationship between  $V \leftrightarrow g$

In spherical harmonics, for  $r > a$

$$V(r, \theta, \phi) = \frac{GM}{r} \left[ 1 + \sum_{l=2}^{\infty} \left( \frac{a}{r} \right)^l \sum_{m=0}^l P_l^m(\cos\theta) (C_l^m \cos m\phi + S_l^m \sin m\phi) \right]$$

$\uparrow$        $\uparrow$        $\uparrow$   
 $V(r)$      $P_l^m$     why no  $l=1$ ?

$$\tilde{g} = \nabla V$$

$$g_r = -\frac{GM}{r^2} \left[ 1 + \sum_{l=2}^{\infty} (l+1) \left( \frac{a}{r} \right)^l \sum_{m=0}^l P_l^m(\cos\theta) (C_l^m \cos m\phi + S_l^m \sin m\phi) \right]$$

$$g_\theta = \frac{1}{r} \frac{\partial V}{\partial \theta} = \frac{GM}{r^2} \sum_{l=2}^{\infty} \left( \frac{a}{r} \right)^l \sum_{m=0}^l \frac{\partial P_l^m}{\partial \theta} (C_l^m \cos m\phi + S_l^m \sin m\phi)$$

$$g_\phi = \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \rightarrow m (-S_l^m \sin m\phi + C_l^m \cos m\phi)$$

For  $\infty$  series,  $g \leftrightarrow V$  provide same information!

Usually look at  $g_r$

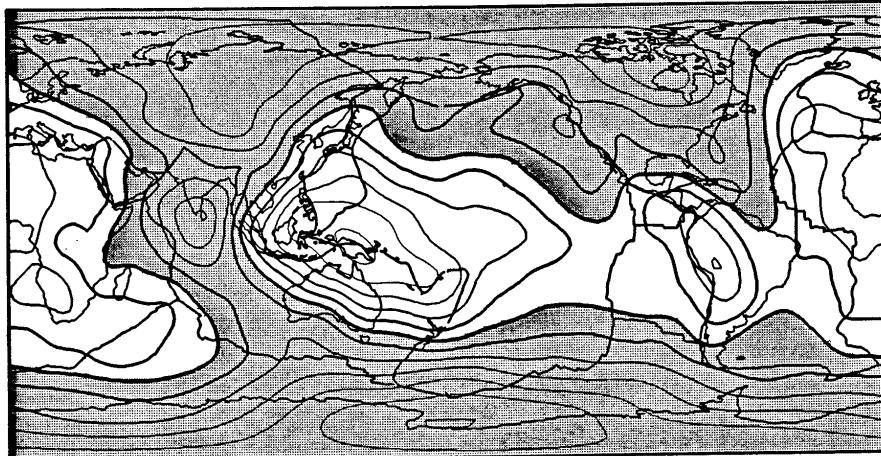
$$g_e^m \sim (l+1) V_e^m \Rightarrow g \text{ emphasizes}$$

high frequency

$V$  emphasizes low frequency

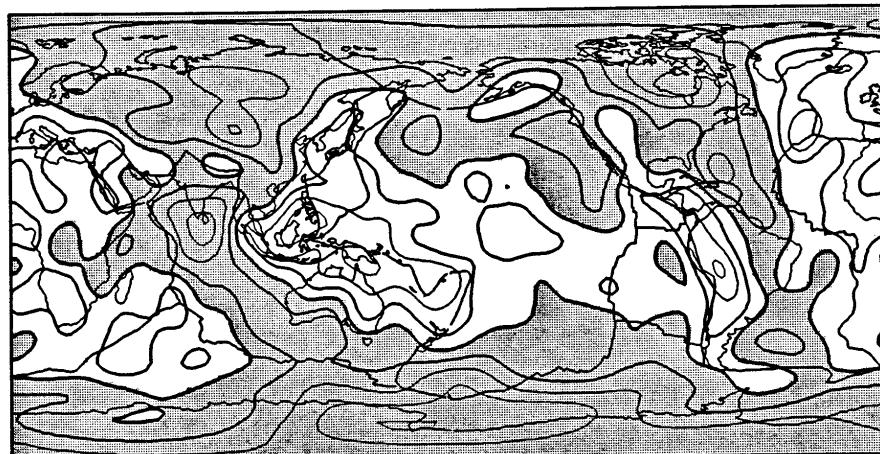
For perfect coverage, both contain same info.

Observed Gravity: degree 6-60



$$S_e \sim l^{-2}$$

contour interval: 20 m  
Free Air Gravity: degree 2-20



$$S_e \sim l^{-1}$$

contour interval: 20 mgal  
Gravity Gradient: degree 2-20



$$S_e \sim l^0$$